



Conventional and gauge couplings with the Δ field

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Motivation



Why to put attention on the isobar $\Delta(1232)\text{MeV}$?

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- ❖ theory?
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- ❖ Spin 3/2 field
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- ❖ Constraints and degrees of freedom
- ❖ Lagrangian and contact transformations
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- ❖ Is this situation exclusive of the RS field?
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- ❖ Conventional and Gauge couplings
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- It is the first excited state of the nucleon and has a prominent role in strong interactions physics.
- Δ -resonance dominates many nuclear phenomena at energies above π -production threshold.



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What we know about the Δ ?

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What we know about the Δ ?

- Was discovered 50 years ago by Fermi and Collaborators at Chicago cyclotron (now Fermilab).

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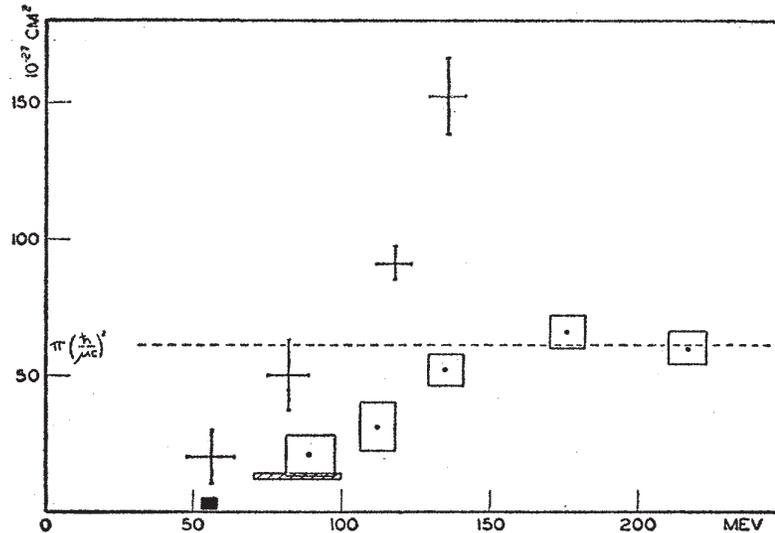


FIG. 1. Total cross sections of negative pions in hydrogen (sides of the rectangle represent the error) and positive pions in hydrogen (arms of the cross represent the error). The cross-hatched rectangle is the Columbia result. The black square is the Brookhaven result and does not include the charge exchange contribution.

- $m = 1232 \text{ MeV} \simeq m_N + 300 \text{ MeV}$, $\Gamma = 120 \text{ MeV}$
($\tau = 10^{-23} \text{ sec}$), $S = T = 3/2$.

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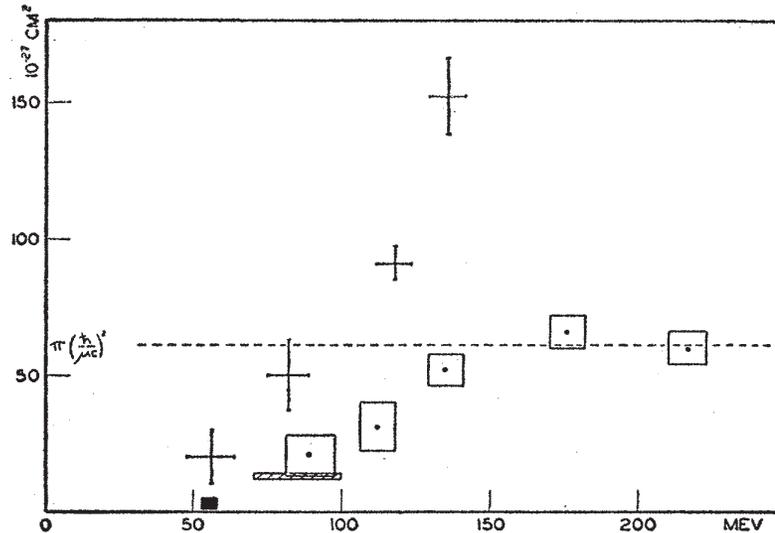


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- In the SU(6) quark model N appears as GS of a 3-q system in a confining potential.



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$$\pi(u\bar{d})p(uyd) \rightarrow \Delta^{++}(uuu) \rightarrow \pi'(u\bar{d})p'(uud)$$



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- Decays 99% of time in πN channel and 1% in the electromagnetic one.



- The Δ electromagnetic excitation is essentially a spin flip magnetic dipole (1) one.

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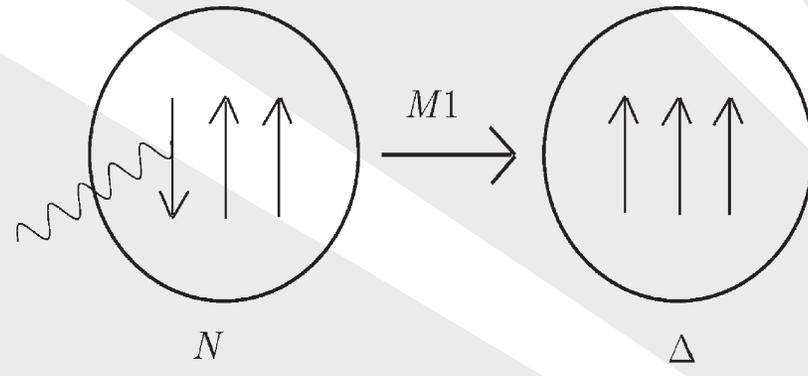
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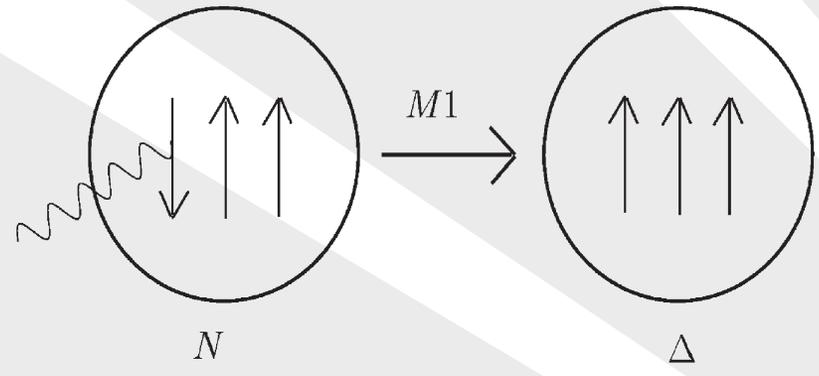
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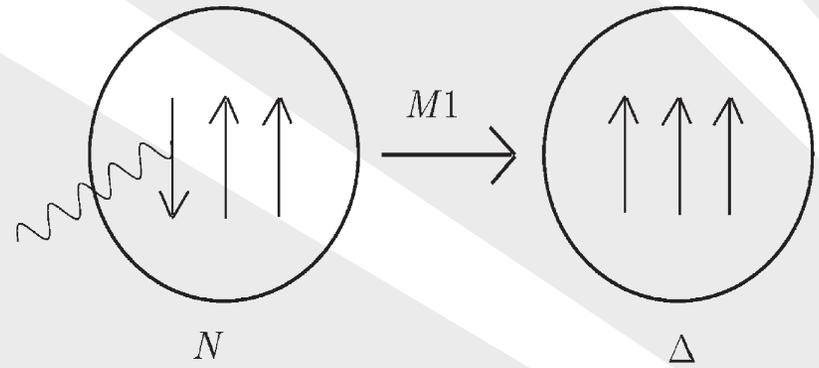
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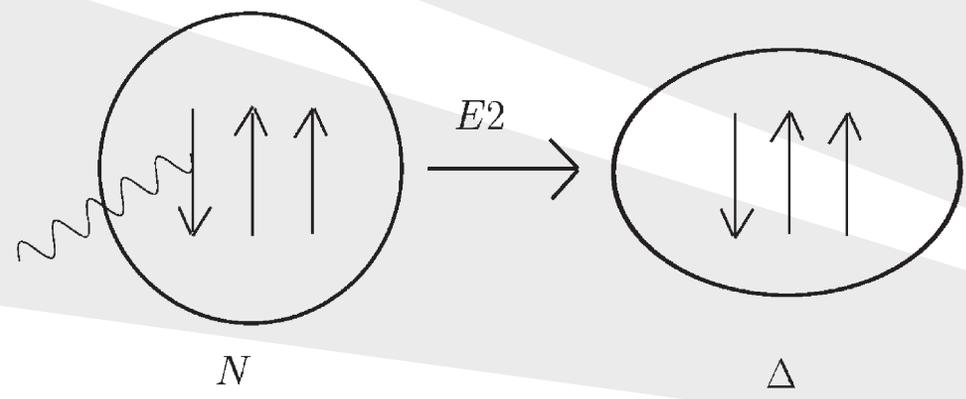
- Or if we considered D admixtures in the N or Δ a quadrupole transition.

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- π are the Goldstone bosons of the spontaneously broken chiral symmetry of QCD + coupling of them
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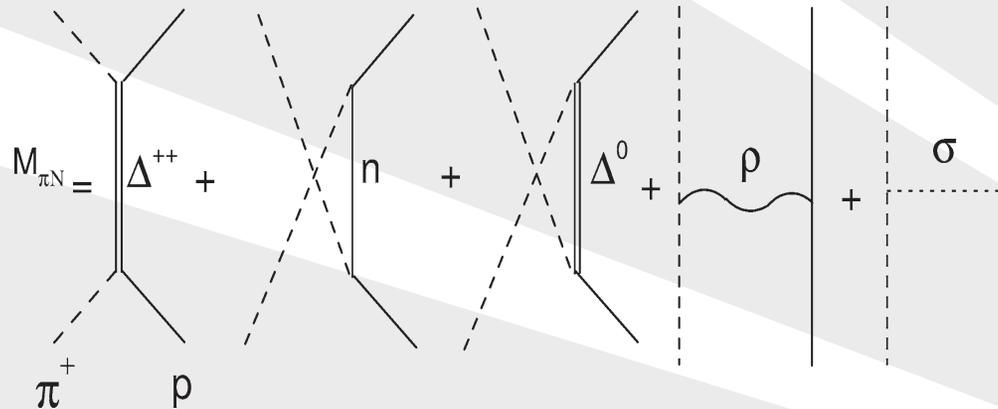
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- Nevertheless pion cloud(.....), effects could only be comprehensively studied within dynamical models, based in the T-matrix calculation.

Spin 3/2 field



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- Rarita Schwinger (RS) spinor ψ_μ is taken as an element of the NU representation of the LG

$$[(1/2, 0) \oplus (0, 1/2)] \otimes [(1/2, 1/2) = (1/2, 0) \otimes (0, 1/2)]$$

$$\blacklozenge D(1/2) \otimes D(0) = D(1/2)$$

$$\blacklozenge D(1/2) \otimes D(1) = D(3/2) \oplus D(1/2),$$

$$\blacklozenge 2(S = 1/2) + 4(3/2) + 2(1/2) + \text{antiparticles} = 16 \text{ dof.}$$

Spin 3/2 field



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- ❖ $\psi_\mu \equiv \psi \otimes W_\mu,$

- ❖ $D(1/2) \otimes D(0) = D(1/2)$

- ❖ $D(1/2) \otimes D(1) = D(3/2) \oplus D(1/2),$

- ❖ $2(S = 1/2) + 4(3/2) + 2(1/2) + \text{antiparticles} = 16 \text{ dof.}$

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- Rarita Schwinger (RS) spinor ψ_μ is taken as an element of the NU representation of the LG

$$[(1/2, 0) \oplus (0, 1/2)] \otimes [(1/2, 1/2) = (1/2, 0) \otimes (0, 1/2)]$$

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- ❖ $2(S = 1/2) + 4(3/2) + 2(1/2) + \text{antiparticles} = 16 \text{ dof.}$

Spin 3/2 field



❖ Motivation

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❖ Spin 3/2 field

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❖ Constraints and degrees of freedom

❖ Lagrangian and contact transformations

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● What equations of motion are satisfied ?

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❖ $\psi_\mu(x)$ satisfies the Klein Gordon eq.

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❖ Contractions $\partial^\mu\psi_\mu$ and $\gamma_5\gamma^\mu\psi_\mu$ both are Dirac spinors

$$(i\partial - m)(\partial^\mu\psi_\mu(x), \gamma_5\gamma^\mu\psi_\mu(x)) = 0. \tag{3}$$

❖ Some of the 16 $\psi_\mu(x)$ ($W_\mu(S = 1)$) satisfy Proca eq.

$$\partial^\nu(\partial_\nu\psi_\mu(x) - \partial_\mu\psi_\nu(x)) + m^2\psi_\mu(x) = 0, \tag{4}$$



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Constraints and degrees of freedom



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Constraints and degrees of freedom



RS field \rightarrow 'constrained' dynamical system \rightarrow supplemented by constraints or subsidiary conditions

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- From the 16 ($4 \otimes 4$) states only 8 satisfy the subsidiary conditions

$$\partial^\mu \psi_\mu = \gamma^\mu \psi_\mu = 0, \quad (5)$$

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- These 8 '3/2 fields' represent physical on-shell $\Delta(1232)$ states and the remaining 8 are 1/2 ones.

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Lagrangian and contact transformations



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Lagrangian and contact transformations



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$$(i\partial - m)\psi_\mu(x) = 0 \quad + \quad \partial^\mu \psi_\mu(x) = \gamma^\mu \psi_\mu(x) = 0, \quad (6)$$

Lagrangian and contact transformations



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Euler Lagrange

$$\mathcal{L}_{free} = \bar{\psi}_\mu(x) \left\{ i\partial_\alpha \Gamma_{\mu\nu}^\alpha - mB_{\mu\nu} \right\} \psi^\nu(x) \quad (\text{RS 1941}), (7)$$

Lagrangian and contact transformations



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where

$$\begin{aligned} \Gamma_{\nu\mu}^\alpha &= g_{\mu\nu} \gamma^\alpha + \frac{1}{3} \gamma_\mu \gamma^\alpha \gamma_\nu - \frac{1}{3} (\gamma_\mu g_\nu^\alpha + g_\mu^\alpha \gamma_\nu), \\ B_{\nu\mu} &= g_{\mu\nu} - \frac{1}{3} \gamma_\mu \gamma_\nu, \end{aligned} \quad (8)$$

Lagrangian and contact transformations



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$$B_{\nu\mu} = g_{\mu\nu} - \frac{1}{3} \gamma_\mu \gamma_\nu, \quad (8)$$

\mathcal{L}_{free} includes constraints and Γ, B do not mix 3/2 with 1/2 states \Rightarrow **only fix 3/2 component of $\psi_\mu(x)$.**



● $\psi_{3/2\mu}$ satisfy Dirac eq. and $\partial^\mu \psi_{3/2\mu} = \gamma^\mu \psi_{3/2\mu} = 0$

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only affects $\psi_{1/2\mu}$ components of ψ_μ and let \mathcal{L}_{free} invariant
 \Rightarrow a whole family of valid one parameter Lagrangians.



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- $\psi_{3/2\mu}$ satisfy Dirac eq. and $\partial^\mu \psi_{3/2\mu} = \gamma^\mu \psi_{3/2\mu} = 0$

- $\psi_{1/2\mu}$ satisfy Dirac eq. and $\partial^\mu \psi_{1/2\mu} \neq 0, \gamma^\mu \psi_{1/2\mu} \neq 0$

- Then contact transformation

$$\psi^\mu \rightarrow \psi'^\mu = R(a)^{\mu\nu} \psi_\nu \equiv (g^{\mu\nu} + a \gamma^\mu \gamma^\nu) \psi_\nu, \quad (9)$$

only affects $\psi_{1/2\mu}$ components of ψ_μ and let \mathcal{L}_{free} invariant
 \Rightarrow a whole family of valid one parameter Lagrangians.

- Put a in terms of a parameter A as $a = 1/2(1 + 3A)$ and make $\psi_\mu \rightarrow \psi'_\mu$ we get (indices omitted)



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● Put a in terms of a parameter A as $a = 1/2(1 + 3A)$ and make $\psi_\mu \rightarrow \psi'_\mu$ we get (indices omitted)

$$\mathcal{L}_{free}(A) = \bar{\psi}(x) R(A) \{i\partial_\alpha \Gamma^\alpha - mB\} R(A) \psi(x), \quad (10)$$

for $A = -1/3 \rightarrow$ RS.



● Finally, $\mathcal{L}_{free}(A)$ is invariant under the change

$$\psi \rightarrow \psi' = R(a)\psi, \quad A \rightarrow A' = \frac{A - 2a}{1 + 4a} \quad a \neq -1/4, A \neq -1/2, (11)$$

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$$G(p, A)_{\mu}^{\beta} (R(A) \{p_{\alpha} \Gamma^{\alpha} + mB\} R(A))_{\beta\nu}^{-1} = g_{\mu\nu},$$



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where

$$G \left(p, -\frac{1}{3} \right)_{\mu\nu} = - \left[\frac{\not{p} + m}{p^2 - m^2} \hat{P}_{\mu\nu}^{3/2} + \frac{2}{m^2} (\not{p} + m) (\hat{P}_{11}^{1/2})_{\mu\nu} + \frac{\sqrt{3}}{m} (\hat{P}_{12}^{1/2} + \hat{P}_{21}^{1/2})_{\mu\nu} \right]. \quad (13)$$

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At difference of the on-shell case where the subsidiary conditions select only the $3/2$ states, when Δ is off-shell ($p^2 \neq m^2$) the $1/2$ ones appear.

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Is this situation exclusive of the RS field?



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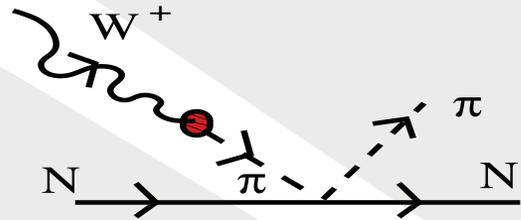
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Is this situation exclusive of the RS field?



W boson in $\pi \rightarrow W \rightarrow \bar{\nu}\mu$ or pion-pole terms in $\nu N \rightarrow \mu N' \pi$



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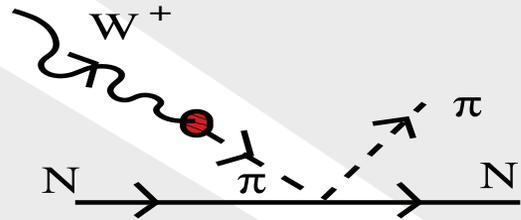
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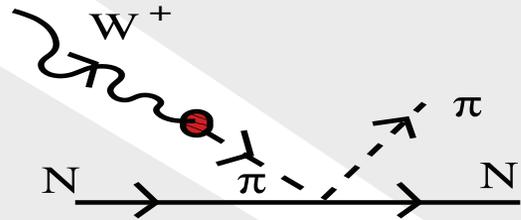
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W field φ_μ is in spin 1 sector of $(1/2, 1/2)$, satisfying subsidiary condition $p^\mu \varphi_\nu = 0$ and the Proca eq.

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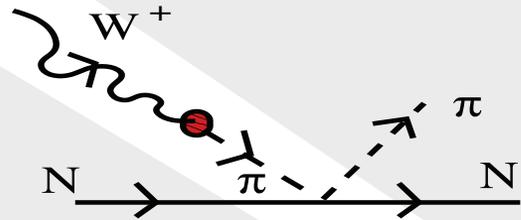
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or equivalently Klein-Gordon one

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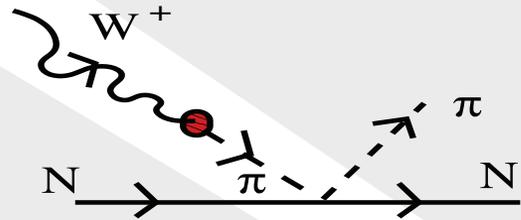
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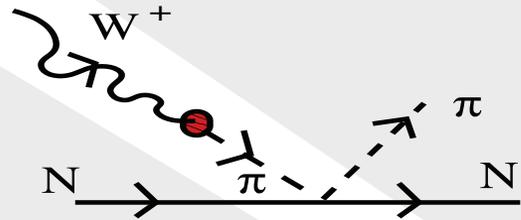
$$(-p^2 + m^2) \varphi_\nu = 0. \quad (15)$$

There is a **spin 0** state satisfying it but $p^\mu \varphi_\nu \neq 0$,

Is this situation exclusive of the RS field?



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or equivalently Klein-Gordon one

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There is a **spin 0** state satisfying it but $p^\mu \varphi_\nu \neq 0$, this is totally analogous to the spin 1/2 and the Dirac eq.

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W propagator looks

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W propagator looks

$$\Delta^{\mu\nu}(p) = - \left[\frac{P_1^{\mu\nu}(p)}{p^2 - m^2} + \frac{P_0^{\mu\nu}(p)}{m^2} \right], \tag{16}$$

being P_0 and P_1 the projectors on the 0 and 1 sectors, respectively.



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Again here an off-shell lower spin contribution as for RS.

Interactions



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Interactions



\mathcal{L}_{int} invariant under contact transformations and leads to A -independent amplitudes.

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Interactions



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We analyze a $\pi(\phi)N(\Psi)\Delta(\psi_\mu)$ vertex,

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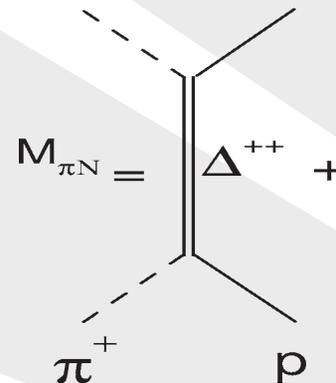
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but can be extended to another interactions.

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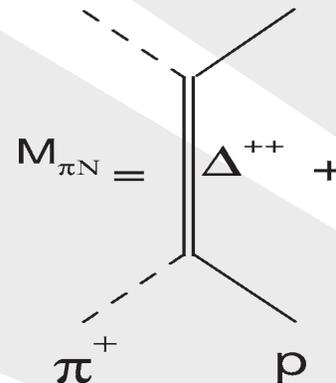
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Interactions



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but can be extended to another interactions. Then (iso. omit.)

$$\mathcal{L}_{int}(A) = g\bar{\psi}_\mu R(A)^{\mu\nu} F_\nu(\psi, \Psi, \phi, \dots) + h.c., \quad (17)$$

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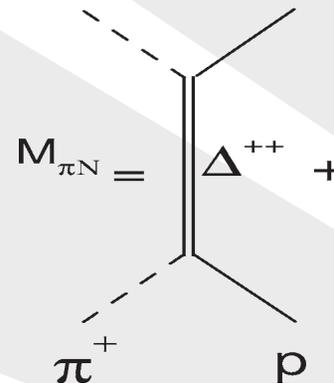
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Interactions



\mathcal{L}_{int} invariant under contact transformations and leads to A -independent amplitudes.

We analyze a $\pi(\phi)N(\Psi)\Delta(\psi_\mu)$ vertex, as appears in



but can be extended to another interactions. Then (iso. omit.)

$$\mathcal{L}_{int}(A) = g\bar{\psi}_\mu R(A)^{\mu\nu} F_\nu(\psi, \Psi, \phi, \dots) + h.c., \quad (17)$$

note $R(A)$ cancel $R(A)^{-1}$ in $G(p, A)$.

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- ❖ **Interactions**
- ❖ Conventional and Gauge couplings

Conventional and Gauge couplings



Two different models for F^μ : the 'conventional' (C) and 'gauge invariant' (G) couplings.

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Nonlinear realization of the chiral symmetry \rightarrow derivative of ϕ

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● $V(p, p_N, p_\pi, -1/3) = -\frac{f_{\pi N \Delta}}{m_\pi} p_\pi^\alpha$

Conventional and Gauge couplings



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Conventional and Gauge couplings



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- Pursue constrain analysis to generate new subsidiary conditions \rightarrow view as dof inconsistency.
- $W \leftrightarrow \pi$ vertex goes as p_π^μ or p_W^μ and $p_{\pi, W}^\mu P_{0\mu\nu} \neq 0 \rightarrow$ impossible for pion decay without coupling the off-shell spin 0 piece of the W propagator.



(G):

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From ($A = -1$ by convenience)

$$\mathcal{L}_{free} = \bar{\psi}_\mu(x) \left(\epsilon^{\mu\nu\alpha\beta} \frac{\partial}{\partial x^\alpha} \gamma_\beta \gamma_5 + im\sigma^{\mu\nu} \right) \psi_\nu(x),$$



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 $F^\mu \partial_\mu \chi = 0,$

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● In massless case $\text{dof} = 2$, (photon case).

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- In massless case $\text{dof} = 2$, (photon case). Mass term breaks this symmetry $\text{dof} = 2 \times 3/2 + 1 = 4$ (part.),

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- To couple a photon $\partial_\mu \rightarrow \partial_\mu - iqA_\mu$



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both gauge symmetries cannot coexist at least order.



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- Evident complexity of the problem,
- C vertex constraints problem not present in perturbative calculations,
- and leaving for a moment the G vertex gauge coexistence,
- only compare C and G in πN elastic scattering within an isobar model.



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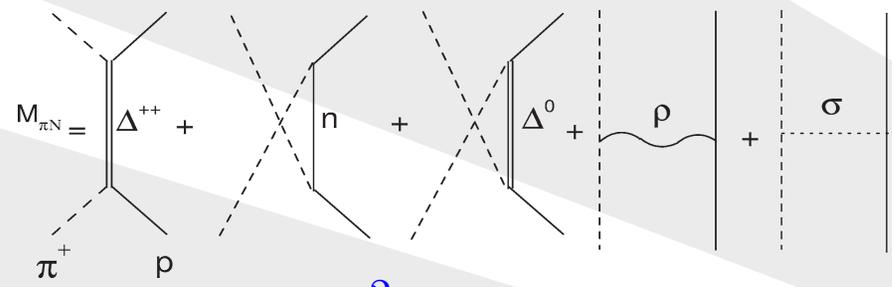
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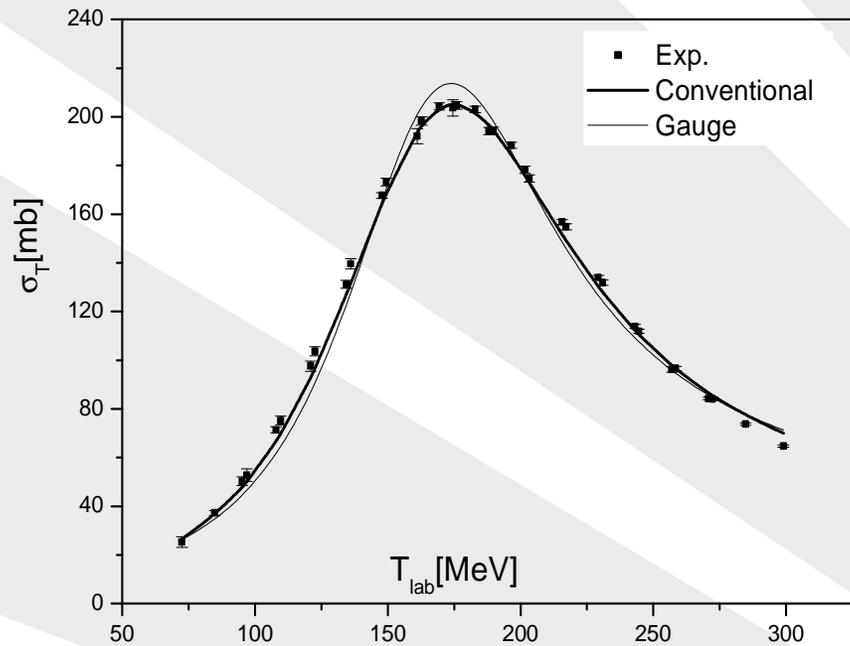


$$\mathcal{M}_{\Delta^{++}}^C = \frac{f_{\pi N \Delta}^2}{m_{\pi}^2} \bar{u}(p'_p, m'_s) p_{\pi}^{\mu} G_{\mu\nu}(p) p_{\pi}^{\nu} u(p_p, m_s),$$

$$\mathcal{M}_{\Delta^{++}}^G = \frac{f_{\pi N \Delta}^2 p^2}{m_{\pi}^2 m^2} \bar{u}(p'_p, m'_s) p_{\pi}^{\mu} (-) \frac{\not{p} + m}{p^2 - m^2} \hat{P}_{\mu\nu}^{\frac{3}{2}} p_{\pi}^{\nu} u(p_p, m_s).$$

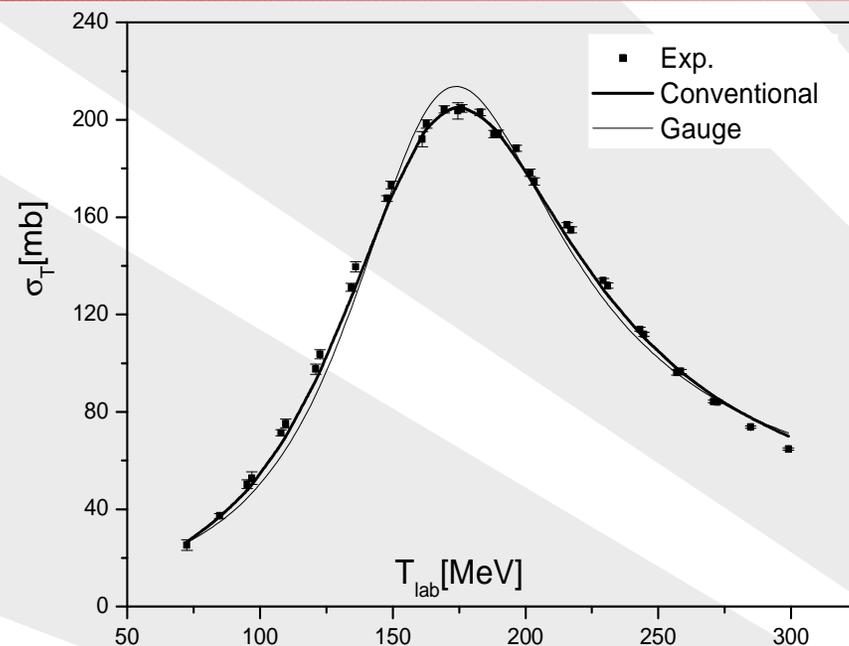
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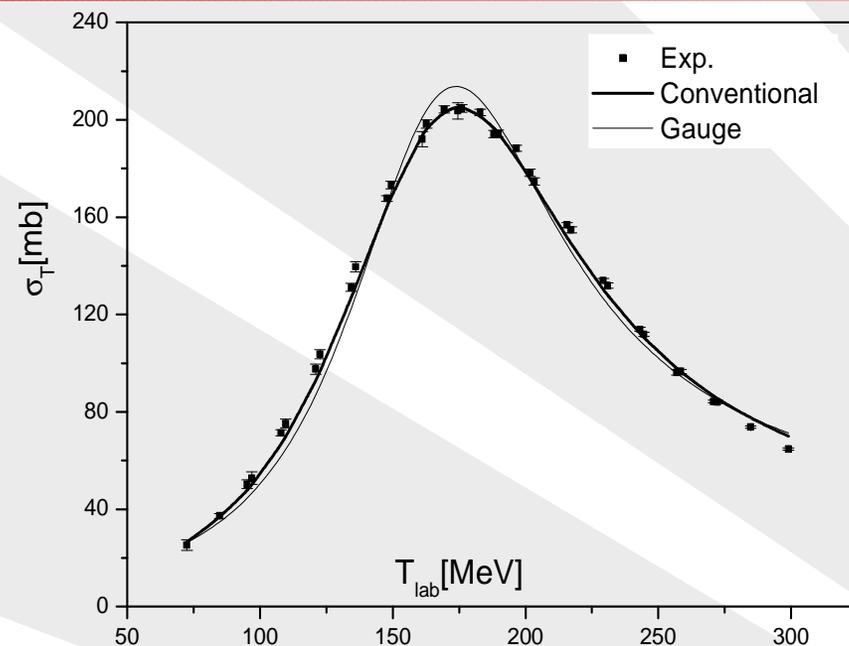
● $m^2 \rightarrow m^2 - im\Gamma$

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- ❖ Constraints and degrees of freedom
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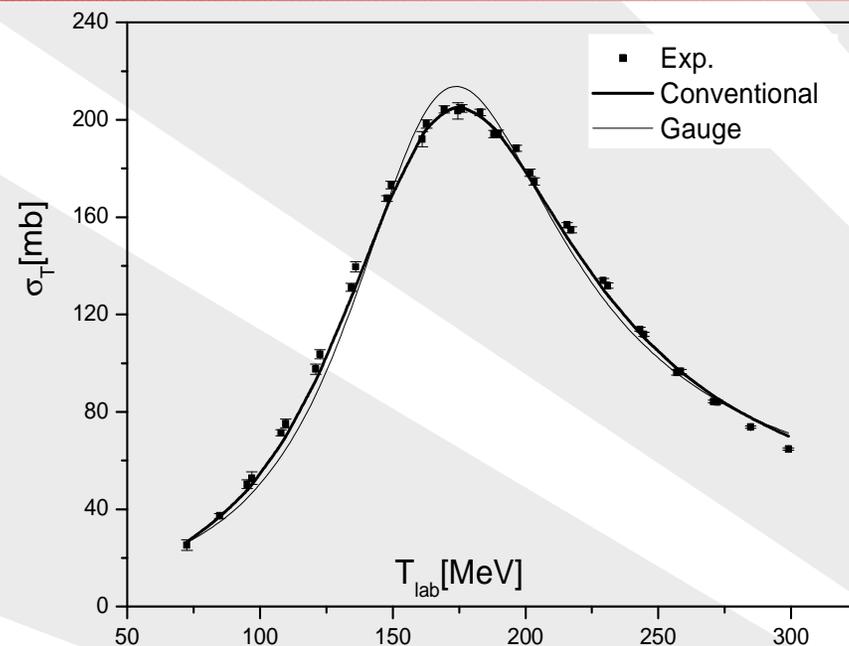
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 $\chi^2/dof = 4.5$

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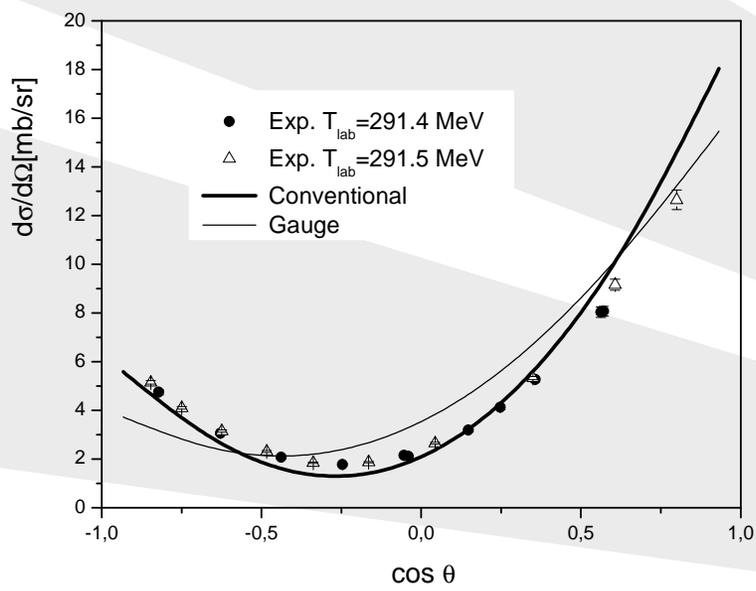
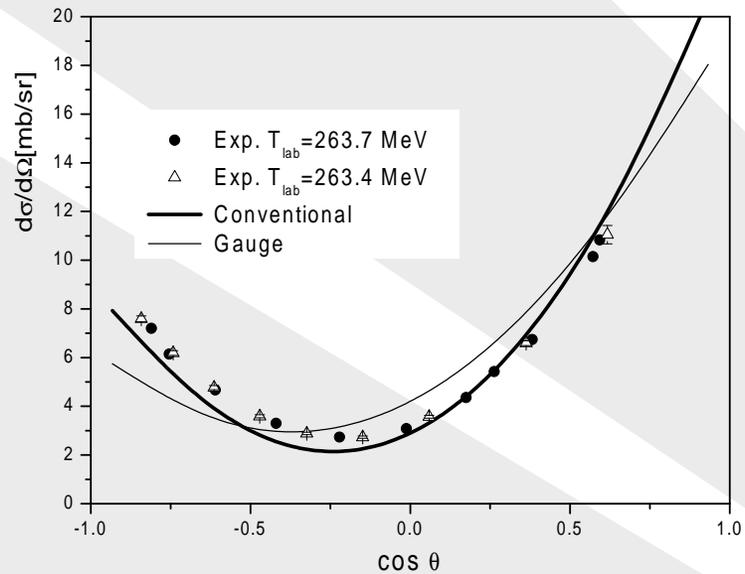
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 1211.6 ± 0.3 MeV, 76.62 ± 0.25 MeV, 1.00 ± 0.05 and 13.5,
 respectively.

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Conclusions



- The so called 'inconsistence' of the off-shell propagation of $1/2$ components is clearly present, in other cases as is the W boson off-shell propagation.

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- The so called 'inconsistence' of the off-shell propagation of $1/2$ components is clearly present, in other cases as is the W boson off-shell propagation.
- Within this simple model, the fitting achieved with the C couplings are clearly better than those obtained with G ones.
- Seems not possible accommodate the parameters of the σ meson (those of ρ are fixed in both approaches by low energy phenomenology) to get identical results with both types of couplings.
- The problem is not closed.