Coherent neutrino scattering

M. Valverde

RCNP - Ōsaka Univ.

45th Karpacz Winter School in Theoretical Physics Ladek-Zdrój, Poland, February 7, 2009

E. Hernandez, J.E. Amaro, J. Nieves and M. V., Phys.Rev.D79:013002,2009

E. Hernandez, J. Nieves and M. V. Phys.Rev.D76:033005,2007

Coherent reaction

$$\left(\overline{
u_I + A_Z|_{gs}}
ightarrow I^- + A_Z|_{gs} + \pi^+
ight)$$

Initial and final nuclear states are the same

- Dominated by Δ mechanism
 Background process largely supressed in isoscalar nuclei
- Important background in neutrino oscillation experiments

The reaction is dominated by nuclear form factor thus it is very forward peaked

Coherent reaction

$$\left(\overline{
u_I + A_Z|_{gs}}
ightarrow I^- + A_Z|_{gs} + \pi^+
ight)$$

Initial and final nuclear states are the same

- Dominated by Δ mechanism
 Background process largely supressed in isoscalar nuclei
- Important background in neutrino oscillation experiments

The reaction is dominated by nuclear form factor thus it is very forward peaked

Outline

1 Free nucleon reaction

2 Nuclear medium effects

Results and comparisons with other models

Outline

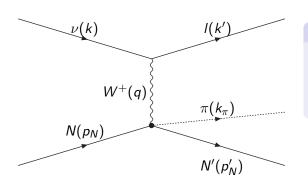
- 1 Free nucleon reaction
- Nuclear medium effects

3 Results and comparisons with other models

Free nucleon reaction

Free nucleon reaction

$$\nu + N \rightarrow I^- + N' + \pi$$



Charge channels

•
$$\nu + p \to I^- + p + \pi^+$$

•
$$\nu + n \to I^- + n + \pi^+$$

•
$$\nu + n \to l^- + n + \pi^+$$

• $(\nu + n \to l^- + p + \pi^0)$

Direct ∆ term

The dominant mechanism is

$$\nu + N \rightarrow I^{-} + \Delta(1232) \rightarrow I^{-} + N' + \pi$$

$$p_{N} \qquad p_{\Delta} \qquad p'_{N}$$

- Resonance dominated studies are most common:
 - Lalakulich and Paschos [PRD71:074003,2005]
 - Graczyk and Sobczyk [PRD77:053001,2008]
- Non-resonant terms have also been studied
 - Sato, Uno and Lee [PRC67:065201,2006]

\triangle -N vertex

$$\begin{split} \left\langle \Delta^{+}; \mathbf{p}_{\Delta} \left| V^{\mu} - A^{\mu} \right| N; \mathbf{p}_{\mathbf{N}} \right\rangle &= \bar{u}_{\alpha}(\mathbf{p}_{\Delta}) \Gamma^{\alpha\mu} u(\mathbf{p}) \cos \theta_{C} \\ \Gamma^{\alpha\mu} &= \left[\frac{C_{3}^{V}}{M} \left(g^{\alpha\mu} \not q - q^{\alpha} \gamma^{\mu} \right) + \frac{C_{4}^{V}}{M^{2}} \left(g^{\alpha\mu} q \cdot p_{\Delta} - q^{\alpha} p_{\Delta}^{\mu} \right) \right. \\ &+ \frac{C_{5}^{V}}{M^{2}} \left(g^{\alpha\mu} q \cdot p - q^{\alpha} p^{\mu} \right) + \frac{C_{6}^{V}}{M^{2}} q^{\alpha} q^{\mu} \right] \gamma_{5} \\ &+ \frac{C_{3}^{A}}{M} \left(g^{\alpha\mu} \not q - g^{\alpha} \gamma^{\mu} \right) + \frac{C_{4}^{A}}{M^{2}} \left(g^{\alpha\mu} q \cdot p_{\Delta} - q^{\alpha} p_{\Delta}^{\mu} \right) \\ &+ \frac{C_{5}^{A}}{M^{2}} g^{\alpha\mu} \right) + \frac{C_{6}^{A}}{M^{2}} q^{\alpha} q^{\mu} \end{split}$$

△-N vertex: form factor parametrization

Vector part

Fit to exp electron scattering data (Lalakulich and Paschos)

$$C_3^V(q^2) = \frac{C_3^V(0)}{(1 - q^2/M_V^2)^2} \frac{1}{1 - q^2/4M_V^2}$$

$$C_5^V = 0, \quad C_4^V = -C_3^V \frac{m_N}{W}, \quad C_3^V(0) = 1.95$$

Only assume CVC

△-N vertex: form factor parametrization

Axial part

Non-diagonal Goldberger-Treiman relation

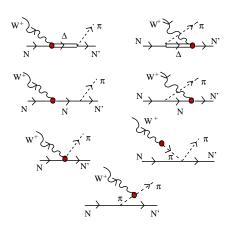
$$C_5^A(q^2) = \frac{C_5^A(0)}{(1 - q^2/M_A^2)^2} \frac{1}{1 - q^2/3M_A^2}$$

$$C_4^A = -\frac{C_5^A}{4}, \quad C_6^A = C_5^A \frac{m_N^2}{m_\pi^2 - q^2}, \quad C_3^A = 0$$

$$C_5^A(0) = \sqrt{\frac{2}{3}} \frac{f_\pi}{m_\pi} f^* = 1.2, \quad M_A = 1.05 \,\text{GeV}$$

Mainly theoretical assumptions

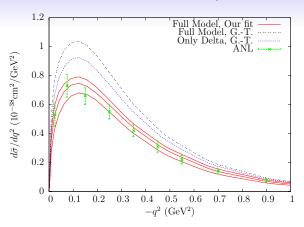
△ and Background terms



- $\Delta(1232)$ resonance pole terms
- Nucleon pole terms
- Other background terms

Non resonant terms computed using chiral SU(2) sigma model with no free parameters

ANL results for $d\sigma/dq^2$



Refit of
$$C_5^A(0)$$

$$C_5^A(0) = 1.2 \rightarrow \boxed{C_5^A(0) = 0.870 \pm 0.075}$$
,

$$M_A = 0.982 \pm 0.082 \text{ GeV}$$

How good is Goldberger-Treiman?

	$C_5^A(0)$			
GT.	1.2			
Our fit	0.870 ± 0.075			
Quark models	0.97	0.83	0.93	0.87
Exp. fit	1.15 ± 0.23	$\boldsymbol{1.39 \pm 0.14}$	1.1 ± 0.2	1.22 ± 0.06
Current algebra	0.98			

Taken from Barquilla-Cano, Buchmann and E. Hernández [PRC75,2007]

Outline

- Free nucleon reaction
- 2 Nuclear medium effects

Results and comparisons with other models

Nuclear Model: Fermi Gas

We describe the nucleus as a Fermi Gas in Local Density Approximation

$$\mathcal{A}^{\mu}_{\pi^+}(q,k_\pi) = \int d^3ec{r} \; \mathrm{e}^{\mathrm{i}\left(ec{q}-ec{k}_\pi
ight)\cdotec{r}}
ho_N(ec{r})\mathcal{J}^{\mu}_{N\pi^+}(ec{r};q,k_\pi)$$

Nuclear effects: Δ is in the medium

Modify
$$M_\Delta o M_\Delta + {\sf Re}\Sigma_\Delta(
ho(r))$$
 and

$$\Gamma_{\Delta}/2
ightarrow \Gamma_{\Delta}^{\mathsf{Pauli}}/2 - \mathsf{Im}\Sigma_{\Delta}(
ho(r))$$

- Potential with atractive real part
- Decay width Γ_{Δ} modified by:
 - Pauli blocking
 - New decay channels are open

[E. Oset and L. L. Salcedo, NPA 468:631,1987]

Outgoing pion FSI

For the pion wave function we solve the Klein-Gordon equation

$$\left[-\vec{\nabla}^2 + \textit{m}_{\pi}^2 + 2\textit{E}_{\pi} \frac{\textit{V}_{\text{opt}}(\vec{r})}{\vec{\varphi}_{\pi^+}^*(\vec{r};\vec{k}_{\pi})} = \textit{E}_{\pi}^2 \widetilde{\varphi}_{\pi^+}^*(\vec{r};\vec{k}_{\pi})\right]$$

Optical imaginary potential

OK because we are dealing with an exclusive reaction

Model by [Nieves, Oset and Garcia-Recio, NPA554,554:1993]

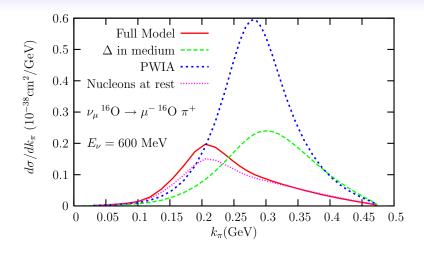
- Δ-dominance at high energies
- Improved model at low energies:
 s and p wave πN in first and second order
- Coulomb effects (Found to be negligible)

Outline

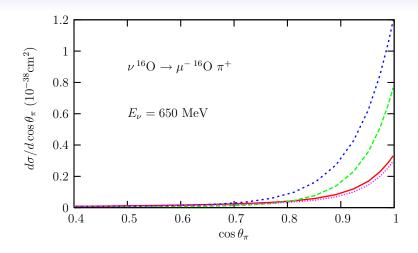
Free nucleon reaction

- 2 Nuclear medium effects
- 3 Results and comparisons with other models

Pion momentum distribution

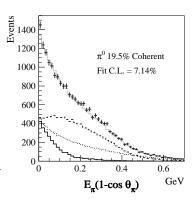


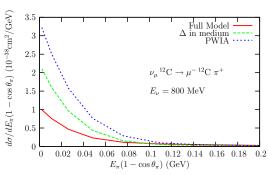
Pion angular distribution



Pion observables for MiniBoone

Nuclear medium effects



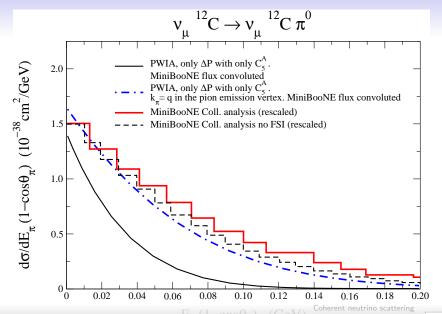


Rein-Sehgal model

Extension of Adler's PCAC theorem

$$\left(\frac{d\sigma_{\nu\nu}}{dxdyd|t|}\right)_{q^2=0} = \left.\frac{G^2ME_{\nu}}{\pi^2}f_{\pi}^2(1-y)|F_{\mathcal{A}}(t)|^2F_{\text{abs}}\frac{d\sigma(\pi^0N\to\pi^0N)}{d|t|}\right|_{q^0=E_{\pi},t=0}$$

It assumes completely forward kinematics so there is no *t*-dependence on the single nucleon matrix element This is good approximation for neutrinos over 2 GeV, but...



Final State Interactions

- In Rein-Sehgal model the pion distortion is taken into account by a Glauber-type absorption factor, but it does not take into account pion absorption or QE scattering
- MiniBoone (NUANCE) uses cascade model but ...
 this is a one step processes (the nuclear final state is known)
 so multiple scattering scattering events should be removed
- We can take into account those QE events in our model by removing that interaction from our optical potential

Integrated cross section results

Reaction	Experiment	$ar{\sigma}~[10^{-40} \mathrm{cm}^2]$	$\sigma_{ m exp}[10^{-40}{ m cm}^2]$
${\sf CC}$ $ u_{\mu}+^{12}{\sf C}$	K2K	4.68	< 7.7
$CC \nu_{\mu} + ^{12}C$	MiniBooNE	2.99	
$CC \nu_{\mu} + ^{12}C$	T2K	2.57	
$CC \nu_{\mu} + ^{16}O$	T2K	3.03	
NC ν_{μ} + 12 C	MiniBooNE	1.97	$7.7\pm1.6\pm3.6$
$NC^* \nu_{\mu} + ^{12}C$	MiniBooNE	2.38*	$7.7\pm1.6\pm3.6$
NC ν_{μ} + 12 C	T2K	1.82	
NC $\nu_{\mu}+^{16}$ O	T2K	2.7	
CC $ar{ u}_{\mu}^{'} + ^{12} C$	T2K	2.12	
NC $\bar{ u}_{\mu}+^{12}$ C	T2K	1.50	

Conclusions

- Coherent is suitable reaction to study $C_5^A(0)$
- Nuclear effects, including pion absorption are crucial
- Rein-Sehgal is not reliable below 2 GeV

Conclusions

- Coherent is suitable reaction to study $C_5^A(0)$
- Nuclear effects, including pion absorption are crucial
- Rein-Sehgal is not reliable below 2 GeV

However, still some work to do

- At high energies other resonances become important
- Exp analysis relies on theoretical analysis which is not so well under control
- Possible nuclear structure effects?