

C5A axial form factor determined from the bubble chamber experiments

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A common project with J. Sobczyk, D. Kiełczewska and P. Przewłocki

Motivations

Nonresonant background negligible

- ∆(1232) excitation induced by vnucleon interaction
- → Simultaneous analysis of the data from two experiments: ANL and BNL
- Extraction of the axial contribution from bubble chamber experiments
 - Fits either to ANL or to BNL data
- Input to →NuWro Monte Carlo Generator
- New fits of cross sections and C5A with account of their uncertainties
- → application to 1π0 production in NC v-nucleon scattering



T. Kitagaki *et al.* Phys. Rev. D 34 (1986) 2554

12 ft @ ANL

- 12 foot bubble chamber filled with deuterium and hydrogen @ Argonne National Laboratory
- S. J. Barish, Phys. Rev. D19 (1979) 2521.
- G. M. Radecky Phys. Rev. D25 (1982) 1161.
- <E> < 1 GeV
- <∆flux> =15% (E<1.5 GeV) and 25% (above)
- Differential cross sections in Q2
- Total cross sections
- Kinematical cuts:
 - 0.5 GeV < E < 6.0 GeV
 - 0.01 GeV2 < Q2 < 1 GeV2
 - W < 1.4 GeV

$$\nu + d \rightarrow \mu^- + \Delta^{++}(1232) + n$$



7-ft @BNL

$$\nu + d \rightarrow \mu^- + \Delta^{++}(1232) + n$$

- 7 foot bubble chamber filled with **deuterium** at Brookhaven National Laboratory.
- <E> = 1.6 GeV
- T. Kitagaki *et al.* Phys. Rev. D 34 (1986) 2554.
- T. Kitagaki *et al.*, Phys. Rev. D42 (1990) 1331.
- $<\Delta flux> = 10\%$
- Kinematical cuts:
 - 0.5 GeV < E < 6.0 GeV
 - Q2 < 3 GeV2 but:</p>
 - $Q2 > 0.1 \rightarrow efficiency!$
 - W<1.4 GeV
- Total cross sections
- Normalized cross sections





Figure 1. Previous measurements of the total cross section per nucleon of the process $\nu_{\mu}p \rightarrow \mu^{-}p\pi^{+}$ at low neutrino energy.

Rarita-Schwinger Formalism

$$\nu + p \rightarrow \mu^- + \Delta^{++} (1232)$$

$$\mathcal{J}_{\mu}^{CC}=\mathcal{J}_{\mu}^{V}+\mathcal{J}_{\mu}^{A},~~$$
Hadronic current

The differential cross section for Δ production is:

$$\frac{d^2\sigma}{dWdQ^2} = \frac{\widetilde{G}^2 W}{64\pi^2 M E^2} L^{\mu\nu} W_{\mu\nu},$$

where $\tilde{G} = G \cos \theta_c$, and G is the Fermi constant, and θ_c Cabibbo angle. The hadronic tensor is:

$$W_{\mu\nu} = \frac{1}{4MM_{\Delta}} \frac{1}{2} \sum_{spin} \left\langle \Delta^{++}, p' \right| \mathcal{J}_{\mu}^{CC} \left| p \right\rangle \left\langle \Delta^{++}, p' \right| \mathcal{J}_{\nu}^{CC} \left| p \right\rangle^* \frac{\Gamma_{\Delta}/2}{((W - M_{\Delta})^2 + \Gamma_{\Delta}^2/4)}$$

The leptonic tensor is:

$$\mathcal{L}_{\mu\nu} = 8\left(k'_{\mu}k_{\nu} + k_{\mu}k'_{\nu} - g_{\mu\nu}k'_{\alpha}k^{\alpha} \mp i\epsilon_{\mu\nu\alpha\beta}k^{\alpha}k'^{\beta}\right).$$

$$\left\langle \Delta^{++}(p') \left| \mathcal{J}^{V}_{\mu} \left| N(p) \right\rangle = \sqrt{3} \bar{\Psi}_{\lambda}(p') \left[g^{\lambda}_{\ \mu} \left(\frac{C_{3}^{V}}{M} \gamma_{\nu} + \frac{C_{4}^{V}}{M^{2}} p'_{\ \nu} + \frac{C_{5}^{V}}{M^{2}} p_{\nu} \right) q^{\nu} - q^{\lambda} \left(\frac{C_{3}^{V}}{M} \gamma_{\mu} + \frac{C_{4}^{V}}{M^{2}} p'_{\mu} + \frac{C_{5}^{V}}{M^{2}} p_{\mu} \right) \right] \gamma_{5} u(p) \right\rangle dp'' - q^{\lambda} \left(\frac{C_{3}^{V}}{M} \gamma_{\mu} + \frac{C_{4}^{V}}{M^{2}} p'_{\mu} + \frac{C_{5}^{V}}{M^{2}} p_{\mu} \right) \right] \gamma_{5} u(p)$$

$$C_5^V(Q^2)=0, \quad C_4^V(Q^2)=-\frac{M}{W}C_3^V(Q^2)$$
 SU(6) symmetry relation

→We apply fits of CiV proposed by O. Lalakulich *et al.* Phys. Rev. D 74 (2006) 014009

$$\left\langle \Delta^{++}(p') \right| \mathcal{J}_{\mu}^{A} \left| N(p) \right\rangle = \sqrt{3} \bar{\Psi}_{\lambda}(p') \left[g^{\lambda}_{\ \mu} \left(\gamma_{\nu} \frac{C_{3}^{A}}{M} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \nu} \right) q^{\nu} - q^{\lambda} \left(\frac{C_{3}^{A}}{M} \gamma_{\mu} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \mu} \right) + g^{\lambda}_{\ \mu} C_{5}^{A} + \frac{q^{\lambda} q_{\mu}}{M^{2}} C_{6}^{A} \right] u(p) \left(q^{\lambda}_{\ \mu} \left(\gamma_{\nu} \frac{C_{3}^{A}}{M} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \nu} \right) q^{\nu} - q^{\lambda} \left(\frac{C_{3}^{A}}{M} \gamma_{\mu} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \mu} \right) + g^{\lambda}_{\ \mu} C_{5}^{A} + \frac{q^{\lambda} q_{\mu}}{M^{2}} C_{6}^{A} \right) u(p) \left(q^{\lambda}_{\ \mu} \left(\gamma_{\nu} \frac{C_{3}^{A}}{M} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \nu} \right) q^{\nu} - q^{\lambda} \left(\frac{C_{3}^{A}}{M} \gamma_{\mu} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \mu} \right) + g^{\lambda}_{\ \mu} C_{5}^{A} + \frac{q^{\lambda} q_{\mu}}{M^{2}} C_{6}^{A} \right) u(p) \left(q^{\lambda}_{\ \mu} \left(\gamma_{\nu} \frac{C_{3}^{A}}{M} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \nu} \right) q^{\nu} - q^{\lambda} \left(\frac{C_{3}^{A}}{M} \gamma_{\mu} + \frac{C_{4}^{A}}{M^{2}} p'_{\ \mu} \right) \right) \right) \right) q^{\mu} dp$$



$$C_5^A(Q^2) = C_5^A(0) \left(1 + \frac{Q^2}{M_A^2}\right)^{-2}$$

Dipole parameterization

From Adler model: **a=-1.21, b= 2 GeV2**

Adler parameterization

$$C_5^A(Q^2) = C_5^A(0) \left(1 + \frac{aQ^2}{b + Q^2}\right) \left(1 + \frac{Q^2}{M_A^2}\right)^{-2}$$

P. A. Schreiner and F. Von Hippel, Nucl. Phys. B 58 (1973) 333.

$$C_5^A(0) = \frac{g_{\pi N\Delta} f_\pi}{\sqrt{6}M} \approx 1.15 \pm 0.01$$

From PCAC see: D. Barquilla-Cano, *et al.* Phys. Rev. C75 (2007) 065203 [Erratum-ibid. C 77 (2008) 019903]

other parameterizations!

$$C_5^A(Q^2) = C_5^A(0) \left(1 + \frac{Q^2}{M_{A1}^2}\right)^{-2} \left(1 + \frac{Q^2}{M_{A2}^2}\right)^{-i=1,2}$$

Adler with a, and b

Parameters seem to be correlated



FIG. 1. Differential cross section for weak charged current neutrino production of Δ on deuteron. In the short-dashed line, deuteron effects are neglected while dotted, long-dashed, and solid lines include these effect using Hulthen, Bonn and Paris deuteron wave functions, respectively.

L.Alvarez-Ruso, S.K.Singh and M.J.Vicente Vacas, Phys. Rev. C 59 (1999) 3386 NN potentials

- Hulthen, L. Hulthen and M.
 Sugawara, Handbuch der Physik
- Bonn, R. Machleidt, K. Holinde and C. Elster, Phys. Rept. 149 (1987) 1
- Paris: M. Lacombe, B. Loiseau, R. Vinh Mau, J. Cote, P. Pires and R. de Tourreil, Phys. Lett. B 101 (1981) 139

Used in this analysis

$$R_{ANL}(Q^2) = \frac{\left(d\sigma(\nu d \to \mu^- n\Delta^{++})/dQ^2\right)_{deuterium}}{(d\sigma(\nu p \to \mu^- \Delta^{++})/dQ^2)_{free\ target}}$$

hi-2 method

$$\chi_{ANL}^{2} = \sum_{i=1}^{n_{ANL}} \left(\frac{\sigma_{th}^{ANL}(Q_{i}^{2}) - p\sigma_{ex}^{ANL}(Q_{i}^{2})}{p\Delta\sigma_{i}^{ANL}} \right)^{2} + \left(\frac{p-1}{r_{ANL}} \right)^{2}$$

$$Total \# sections$$
Flux uncertainty
$$\sigma_{tot-ex}^{ANL} = \sum_{i=1}^{n_{ANL}} \Delta Q_{i}^{2} \sigma_{ex}^{ANL}(Q_{i}^{2}) = 0.31278 \times 10^{-38} \text{ cm}^{2},$$

Analogically the $\sigma_{\textit{th}}$ cross section is obtained

$$\sigma_{th}^{ANL}(Q_i^2) = \frac{1}{\Psi_{ANL}} \cdot \frac{1}{\Delta Q_i^2} \int_{Q_i^2 - \Delta Q_i^2/2}^{Q_i^2 + \Delta Q_i^2/2} dQ^2 \int_{E_{min}}^{E_{max}} dE \Phi_{ANL}(E) \sigma_{th}(Q^2, E), \tag{38}$$

where the $\Phi_{ANL}(E)$ is the neutrino flux, which is normalized to the area restricted by $E_{min} = 0.5$ GeV and $E_{max} = 6.0$ GeV.

$$\Psi_{ANL} = \int_{E_{min}}^{E_{max}} dE \Phi_{ANL}(E) \tag{39}$$

 $\chi^{2} = \chi^{2}_{ANL} + \chi^{2}_{BNL} \qquad \qquad n_{ANL} + n_{BNL} + 1 - n_{par} - 2$

$C_5^A(Q^2) = C_5^A(0) \bigg($	$1 + \frac{Q^2}{M_A^2} \bigg)^{-2}$	$M_A \; [\mathrm{GeV}^2]$	$C_5^A(0)$	PANL	<i>p</i> bnl	χ^2/NDF	GoF	Ρ
	free target	0.96 ± 0.04	1.15(fixed)	1.16 ± 0.06	0.98 ± 0.03	0.89	0.62	R
		0.95 ± 0.04	1.14 ± 0.08	1.15 ± 0.10	0.98 ± 0.03	0.92	0.57	F
	deuteron	0.94 ± 0.04	1.15(fixed)	1.04 ± 0.06	0.97 ± 0.03	0.87	0.66	Ъ.
		0.94 ± 0.03	1.19 ± 0.09	1.08 ± 0.10	0.98 ± 0.03	0.88	0.64	
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$C_5^A(Q^2) = C_5^A(0) \left(1 + \frac{aQ^2}{b+Q} \right)$	$\frac{1}{2}\left(1+\frac{Q^2}{M_A^2}\right)^{-2}$	$M_A \ [\text{GeV}^2]$	$C_5^A(0)$	p_{ANL}	p_{BNL}	χ^2/NDF	GoF
	free target	1.32 ± 0.07	1.15(fixed)	1.24 ± 0.06	0.98 ± 0.03	0.94	0.55
For the BNL Q2 > 0.1!		1.32 ± 0.07	1.09 ± 0.08	1.18 ± 0.10	0.98 ± 0.03	0.96	0.51
	deuteron	1.29 ± 0.07	1.15(fixed)	1.12 ± 0.06	0.98 ± 0.03	0.86	0.67
a=-1.21, b= 2	GeV2	1.29 ± 0.07	1.13 ± 0.08	1.10 ± 0.10	0.98 ± 0.03	0.89	0.62

$C_5^A(Q^2) = C_5^A(0) \left(1 + \frac{Q^2}{M_A^2}\right)^{-2}$	$M_A \; [{ m GeV}^2]$	$C_{5}^{A}(0)$	p_{ANL}	p_{BNL}	χ^2/NDF	GoF
free target	0.96 ± 0.04	1.15(fixed)	1.16 ± 0.06	0.98 ± 0.03	0.89	0.62
	0.95 ± 0.04	1.14 ± 0.08	1.15 ± 0.10	0.98 ± 0.03	0.92	0.57
About 11% of between ANL	differenc and BNL	e 1.15(fixed)	1.04 ± 0.06	0.97 ± 0.03	0.87	0.66
9-10% of deu structure effe	teron cts	1.19 ± 0.09	1.08 ± 0.10	0.98 ± 0.03	0.88	0.64



$C_5^A(Q^2) = C_5^A(0) \left(1 + \frac{Q^2}{M_A^2} \right)$	$\int^{-2} M_A \ [\text{GeV}^2]$	$C_5^A(0)$	PANL	p_{BNL}	χ^2/NDF	GoF
free tar	get 0.96 ± 0.04	1.15(fixed)	1.16 ± 0.06	0.98 ± 0.03	0.89	0.62
	0.95 ± 0.04	1.14 ± 0.08	1.15 ± 0.10	0.98 ± 0.03	0.92	0.57
deuter	$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	1.15(fixed)	1.04 ± 0.06	0.97 ± 0.03	Dipole 0.94 G	axial mass eV
	0.94 ± 0.03	1.19 ± 0.09	1.08 ± 0.10	0.98 ± 0.03	0.88	0.64











Number of events/0.05 GeV²



With deuteron structure effects

Summary

- Fits are self consistent
- The deuteron structure effects must be accounted especially for ANL data
- The relative difference between ANL and BNL data is around 11 % in the normalization
- The C5A(0) value is consistent with PCAC constraint
 - The analysis of uncertainties of cross sections (due to data) for $1\pi0$ production
 - A model independent analysis of the data?.....→ Neural Networks?